

Measurement of the $B^0-\bar{B}^0$ mixing rate with $B^0(\bar{B}^0)\rightarrow D^{*\mp}\pi^\pm$ partial reconstruction

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We report a measurement of the $B^0 - \bar{B}^0$ mixing parameter Δm_d based on a 29.1 fb^{-1} sample of $\Upsilon(4S)$ resonance decays collected by the Belle detector at the KEKB asymmetric e^+e^- collider. We use events with a partially reconstructed $B^0(\bar{B}^0) \rightarrow D^{*\mp} \pi^\pm$ candidate and where the flavor of the accompanying B meson is identified by the charge of the lepton from a $B^0(\bar{B}^0) \rightarrow X^\mp \ell^\pm \nu$ decay. The proper-time difference between the two B mesons is determined from the distance between the two decay vertices. From a simultaneous fit to the proper-time distributions for the same-flavor ($B^0(\bar{B}^0), \ell^\pm$) and opposite-flavor ($B^0(\bar{B}^0), \ell^\mp$) event samples, we measure the mass difference between the two mass eigenstates of the neutral B meson to be $\Delta m_d = (0.509 \pm 0.017 \text{ (stat)} \pm 0.020 \text{ (syst)}) \text{ ps}^{-1}$.

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I. INTRODUCTION

After production, B^0 and \bar{B}^0 mesons evolve in time and mix into each other via the second-order weak interaction box diagrams shown in Fig. 1. The mixing parameter Δm_d , which is the mass difference of the two neutral B mass eigenstates, is determined from the Feynman diagrams shown in Fig. 1 [1] to be

$$\Delta m_d = \frac{G_F^2}{6\pi^2} f_B^2 m_B m_W^2 \eta_t S |V_{tb}^* V_{td}|^2 B_B, \quad (1)$$

where m_t , m_B and m_W are the t -quark, B^0 and W masses; G_F is the Fermi constant; η_t is a QCD correction [2]; S is a function of m_t^2/m_W^2 [3]; f_B is the decay constant of the B meson; and B_B is the B meson bag parameter [4]. In principle, a measurement of the mixing parameter Δm_d can be used to determine the magnitude of the CKM matrix element $|V_{td}|$. However, there are large theoretical uncertainties associated with the model dependence of f_B and B_B . We report here a measurement of Δm_d that uses B mesons from the decays of $\Upsilon(4S)$ states produced by the KEKB collider and recorded in the Belle detector. We determine the flavor of one B meson by partially reconstructing the decays $B^0(\bar{B}^0) \rightarrow D^{*\mp} \pi^\pm$; the flavor of the accompanying B meson is iden-

tified by the charge of the lepton from $B^0(\bar{B}^0) \rightarrow X^\mp \ell^\pm \nu$ decays. In the future, this technique can be extended to determine the linear combination of CKM angles $2\phi_1 + \phi_3$ [5].

The $\Upsilon(4S)$ decays to a $B^0 - \bar{B}^0$ pair that is nearly at rest in the $\Upsilon(4S)$ center of mass system (c.m.). Since B mesons are spin 0 mesons, angular momentum conservation requires the two B mesons to be in an antisymmetric quantum state

$$|\alpha\rangle = \frac{1}{\sqrt{2}} (|B^0(1)\rangle |\bar{B}^0(2)\rangle - |\bar{B}^0(1)\rangle |B^0(2)\rangle), \quad (2)$$

where 1 and 2 indicate the opposite sides in the $\Upsilon(4S)$ decay plane. The form of Eq. (2) guarantees that at any time the amplitude for either $B^0\bar{B}^0$ or $\bar{B}^0 B^0$ states vanishes. Thus, if we determine the flavor and decay time t_{tag} for one of the B mesons to decay into a final state f_{tag} , then we can measure the time evolution of the other B at any time t as a function of the time difference $\Delta t = t - t_{tag}$. For the two-state neutral B system, the following probability expressions hold to a good approximation [1]:

$$\text{P}^{\text{OF}}(\Delta t) = \frac{1}{4\tau_{B^0}} e^{-|\Delta t|/\tau_{B^0}} [1 + \cos(\Delta m_d \Delta t)],$$

$$\text{P}^{\text{SF}}(\Delta t) = \frac{1}{4\tau_{B^0}} e^{-|\Delta t|/\tau_{B^0}} [1 - \cos(\Delta m_d \Delta t)], \quad (3)$$

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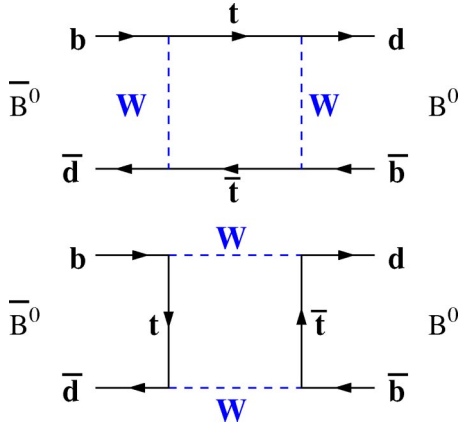


FIG. 1. Standard model “box diagrams” for the second-order weak $B^0-\bar{B}^0$ mixing process.

where the superscripts SF and OF denote events where the lepton-tagged and partially reconstructed B mesons have the same and opposite flavors, respectively.

Since the B^0 and \bar{B}^0 are nearly at rest in the c.m. frame, Δt can be determined from the displacement between the lepton-tagged and partially reconstructed B decay vertices

$$\Delta t \simeq (z_{\text{rec}} - z_{\text{tag}}) / \beta \gamma c \equiv \Delta z / \beta \gamma c, \quad (4)$$

where the z axis is defined to be anti-parallel to the positron beam direction and the constant $\beta \gamma = 0.425$ is the Lorentz boost of the e^+e^- center of mass system at the KEKB collider.

This paper is organized as follows. In Sec. II, we describe the KEKB collider and the Belle detector. The partial reconstruction of $B^0(\bar{B}^0) \rightarrow D^{*\mp} \pi^\pm$ decays, the determination of the b -flavor of the accompanying B meson and the measurement of Δt are described in Sec. III. The likelihood fit to the measured Δt distributions is described in Sec. IV. We present the results of the fit and studies of sources of systematic errors in Secs. V, VI and VII. The conclusions are presented in Sec. VIII.

II. EXPERIMENTAL APPARATUS

KEKB [6] is an asymmetric e^+e^- collider 3 km in circumference, which consists of 8 GeV e^- and 3.5 GeV e^+ storage rings and an injection linear accelerator. It has a single interaction point (IP) where the e^+ and e^- beams collide with a crossing angle of 22 mrad. The collider has reached a peak luminosity above $8 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$. Due to the energy asymmetry, the $Y(4S)$ resonance and its daughter B mesons are produced with a Lorentz boost of $\beta \gamma = 0.425$. On average, the B mesons decay approximately 200 μm from the $Y(4S)$ production point.

The Belle detector [7] is a general-purpose large solid angle magnetic spectrometer surrounding the interaction point. Precision tracking and vertex measurements are provided by a silicon vertex detector (SVD) [8] and a central drift chamber (CDC) [9] in a 1.5 T magnetic field parallel to the z axis. The SVD consists of three layers of double-sided

silicon strip detectors (DSSD) arranged in a barrel and covers 86% of the solid angle. The three layers at radii of 3.0, 4.5 and 6.0 cm surround the beam-pipe, a double-wall beryllium cylinder of 2.3 cm radius and 1 mm thickness. The strip pitches of each DSSD are 84 μm for the measurement of the z coordinate and 25 μm for the measurement of the $r-\phi$ coordinate. The CDC is a small-cell cylindrical drift chamber with 50 layers of anode wires including 18 layers of stereo wires. A low- Z gas mixture [$\text{He}(50\%) + \text{C}_2\text{H}_6(50\%)$] is used to minimize multiple Coulomb scattering to ensure good momentum resolution, especially for low momentum particles. The CDC provides three-dimensional trajectories of charged particles in the polar angle region $17^\circ < \theta < 150^\circ$ in the laboratory frame. The impact parameter resolution for reconstructed tracks is measured as a function of the track momentum p (measured in GeV/c) to be $\sigma_{xy} = [19 \oplus 50 / (p \beta \sin^{3/2} \theta)] \mu\text{m}$ and $\sigma_z = [36 \oplus 42 / (p \beta \sin^{5/2} \theta)] \mu\text{m}$. The momentum resolution of the combined tracking system is $\sigma_{p_t} / p_t = (0.30 / \beta \oplus 0.19 p_t) \%$, where p_t is the transverse momentum in GeV/c .

The identification of charged pions and kaons uses three detector systems: the CDC measuring dE/dx , a set of time-of-flight counters (TOF) [10] and a set of aerogel Cherenkov counters (ACC) [11]. The CDC measures energy loss for charged particles with a resolution of $\sigma(dE/dx) = 6.9\%$ for minimum-ionizing pions. The TOF consists of 128 plastic scintillators viewed on both ends by fine-mesh photo-multipliers that operate stably in the 1.5 T magnetic field. Their time resolution is 95 ps (rms), providing three standard deviation (3σ) K^\pm/π^\pm separation below 1.0 GeV/c , and 2σ separation up to 1.5 GeV/c . The ACC consists of 1188 aerogel blocks with refractive indices between 1.01 and 1.03 depending on the polar angle. Fine-mesh photo-multipliers detect the Cherenkov light. The effective number of photo-electrons is approximately 6 for $\beta=1$ particles. Using the information from these three particle identification systems, the K/π likelihood ratio $P(K/\pi) = \mathcal{L}(K) / (\mathcal{L}(K) + \mathcal{L}(\pi))$ is calculated, where $\mathcal{L}(K)$ and $\mathcal{L}(\pi)$ are kaon and pion likelihoods [7]. A selection with $P(K/\pi) > 0.6$ retains about 90% of the charged kaons with a charged pion misidentification rate of about 6%.

Photons and other neutral particles are reconstructed in a CsI(Tl) crystal calorimeter (ECL) [12] consisting of 8736 crystal blocks, 16.2 radiation lengths (X_0) thick. Electron identification is based on a combination of dE/dx measurements in the CDC, the response of the ACC, the position and the shape of the electromagnetic shower, as well as the ratio of the cluster energy to the particle momentum [13]. For the electron identification requirement used in this analysis, the electron identification efficiency is determined from two-photon $e^+e^- \rightarrow e^+e^-e^+e^-$ processes to be more than 90% for $p_{\text{lab}} > 1.0 \text{ GeV}/c$. The hadron misidentification probability, determined using tagged pions from inclusive $K_S^0 \rightarrow \pi^+\pi^-$ decays, is below 0.5%.

All the detectors mentioned above are inside a superconducting solenoid of 1.7 m radius. The outermost spectrometer subsystem is a K_L^0 and muon detector (KLM) [14], that consists of 14 layers of iron (4.7 cm thick) absorber

alternating with resistive plate counters (RPC). The KLM system covers polar angles between 20 and 155 degrees. The efficiency of the muon identification requirement used here, determined by using the two-photon process $e^+e^- \rightarrow e^+e^-\mu^+\mu^-$ and simulated muons embedded in $B\bar{B}$ candidate events, is greater than 90% for tracks with $p_{\text{lab}} > 1 \text{ GeV}/c$. The corresponding pion misidentification probability, determined using $K_S^0 \rightarrow \pi^+\pi^-$ decays, is less than 2%.

In our analysis, Monte Carlo (MC) events are generated using the QQ event generator [15] and the response of the Belle detector is precisely simulated by a GEANT3-based program [16]. The simulated events are then reconstructed and analyzed with the same procedure as is used for the real data.

III. EVENT RECONSTRUCTION

A. Data sample

We analyze a 29.1 fb^{-1} data sample recorded on the $Y(4S)$ resonance. The data was taken from June 1999 to July 2001 and corresponds to about $3.13 \times 10^7 B\bar{B}$ pairs.

B. $B\bar{B}$ event pre-selection

In Belle, neutral B mesons can only be created via the process $Y(4S) \rightarrow B^0\bar{B}^0$. To suppress the non- $b\bar{b}$ background processes from QED, beam-gas and $e^+e^- \rightarrow \tau^+\tau^-$, we select hadronic events using event multiplicity and total energy variables [17].

C. $B \rightarrow D^*\pi$ decay partial reconstruction

We now describe a partial reconstruction method for the decay chain $B \rightarrow D^*\pi_f$, $D^* \rightarrow D\pi_s$, where π_f and π_s designate a fast π and slow π , respectively.¹ This is a variation of a technique that was first developed by the CLEO Collaboration [18]. For the Belle experiment, we modify and apply this method to make a precise time dependent measurement of Δm_d . Unlike analyses that use fully reconstructed $B \rightarrow D^*\pi_f$ decays, we do not use any properties of the decay products of the D^0 meson in the decay $D^* \rightarrow D\pi_s$. According to a Monte Carlo simulation, this method yields an order of a magnitude more events than a full reconstruction method.

1. Kinematics

We consider five particles in the decay chain: B , D^* , π_f , D and π_s . When reconstructed in this way, the system has 20 degrees of freedom.

For partial reconstruction, only the π_f and π_s are used. The D candidate is not reconstructed. We can obtain constraints from 4-momentum conservation for both the decays $B \rightarrow D^*\pi_f$ and $D^* \rightarrow D\pi_s$ (8 constraints). The B , D^* , D , π_f and π_s masses from the PDG2000 compila-

tion [19] provide further 5 constraints. In addition, we use the constraint that the B energy is the CM beam energy of KEKB at the $Y(4S)$ divided by 2 (1 constraint). The measurements of the π_f and π_s 3-vectors provide 6 constraints. In total there are 20 constraints, equal to the number of degrees of freedom of the system. Following previous analyses [18,20], we use two variables to measure the $B \rightarrow D^*\pi$ signal. The first is D^0 missing mass,

$$M_{D_{\text{miss}}}^2 = m_B^2 + m_{\pi_f}^2 + m_{\pi_s}^2 - 2E_B E_{\pi_f} - 2E_B E_{\pi_s} + 2E_{\pi_f} E_{\pi_s} + 2|\vec{p}_B||\vec{p}_{\pi_f}|\cos\theta_{B\pi_f} + 2|\vec{p}_B||\vec{p}_{\pi_s}|\cos\theta_{B\pi_s} - 2|\vec{p}_{\pi_f}||\vec{p}_{\pi_s}|\cos\theta_{\pi_f\pi_s}, \quad (5)$$

where E , \vec{p} , m are the energy, momentum and nominal masses for the B , D^* , D , π_f and π_s mesons in the c.m. frame; $\theta_{B\pi_{f(s)}}$ is the angle between the directions of motion of the B and $\pi_{f(s)}$; and $\theta_{\pi_f\pi_s}$ is the angle between the π_f and π_s . In the c.m. frame, $\cos\theta_{B\pi_s} \approx -\cos\theta_{B\pi_f}$; the analysis [20] shows that this approximation, and the relation

$$\cos\theta_{B\pi_f} = \frac{-m_B^2 - m_{\pi_f}^2 + m_{D^*}^2 + 2E_B E_{\pi_f}}{2|\vec{p}_B||\vec{p}_{\pi_f}|}, \quad (6)$$

can be used to evaluate $M_{D_{\text{miss}}}$ without reconstructing the B meson's flight direction. The second variable is the angle $\theta_{\pi_s}^*$ between the slow pion in the D^* rest frame and the direction of motion of the D^* in the c.m. frame. Using partial reconstruction, it is calculated using the relation,

$$\cos\theta_{\pi_s}^* = \frac{\beta_{D^*}(E_D^* - E_{\pi_s}^*)}{2|\vec{p}_{\pi_s}^*|} - \frac{|\vec{p}_D|^2 - |\vec{p}_{\pi_s}|^2}{2\gamma_{D^*}^2\beta_{D^*}m_{D^*}|\vec{p}_{\pi_s}^*|}, \quad (7)$$

where $\gamma_{D^*} \equiv E_{D^*}/m_{D^*} = (E_B - \sqrt{|\vec{p}_{\pi_f}|^2 + m_{\pi_f}^2})/m_{D^*}$, $\beta_{D^*} \equiv \sqrt{1 - (1/\gamma_{D^*}^2)}$ and $|\vec{p}_D| = \sqrt{E_D^2 - m_D^2} = \sqrt{(E_B - E_{\pi_f} - E_{\pi_s})^2 - m_D^2}$; E_D^* and $E_{\pi_s}^*$ are the energy of D and π_s in the D^* rest frame; $\vec{p}_{\pi_s}^*$ is the momentum of π_s in the D^* rest frame. The asterisks denote variables calculated in the D^* rest frame. Calculated in this way, $\cos\theta_{\pi_s}^*$ can take values outside the physical region due to finite resolutions, or for background events.

2. $B^0(\bar{B}^0) \rightarrow D^{*\mp}\pi^\pm$ event selection

The $B^0(\bar{B}^0) \rightarrow D^{*\mp}\pi^\pm$ candidates are reconstructed with the following requirements. In the c.m. frame, we select π_f candidates with momentum \vec{p}_{π_f} in the range $2.05 \text{ GeV}/c < |\vec{p}_{\pi_f}| < 2.45 \text{ GeV}/c$ and an oppositely charged π_s with momentum $|\vec{p}_{\pi_s}|$ less than $0.45 \text{ GeV}/c$. We require $dr < 0.05 \text{ cm}$, $dz < 2 \text{ cm}$ for the π_f candidate and $dr < 0.2 \text{ cm}$, $dz < 2 \text{ cm}$ for the π_s candidate to suppress backgrounds from beam particles that interact with the residual

¹Throughout the paper, the charge conjugate process is implied, e.g. $B^0 \rightarrow D^{*-}\pi^+$ denotes also $\bar{B}^0 \rightarrow D^{*+}\pi^-$ etc.

gas of the vacuum system or spent-beam particles that strike the vacuum chamber wall. The variables dr and dz are the distances of closest approach of the track to the interaction point in the $r-\phi$ and z planes, respectively. For better vertex resolution, the SVD is required to provide at least two spatial points for a π_f candidate track. The π_f candidates are required to have electron and muon likelihood ratios less than 0.8 to suppress background from semileptonic B decay. Angular momentum conservation in the pseudoscalar to vector-pseudoscalar decay $B^0 \rightarrow D^{*-} \pi^+$ leads to a distribution proportional to $\cos^2 \theta_{\pi_s}^*$ for the angle $\theta_{\pi_s}^*$. To enhance the signal to background ratio, we require $0.3 < |\cos \theta_{\pi_s}^*| < 1.05$. We then select $B^0 \rightarrow D^{*-} \pi^+$ candidates with D^0 missing mass $M_{D_{miss}}$ greater than $1.85 \text{ GeV}/c^2$ and $0.3 < |\cos \theta_{\pi_s}^*| < 1.05$. These two requirements define the signal region.

Figure 2 shows that the signal peaks at zero in the variable $|\cos \theta_{-\pi_f \pi_s} - \cos \theta_{-\pi_f D^*} \cos \theta_{\pi_s D^*}|$ [20]. For about 4% of the events, more than one pair of opposite sign particles satisfy all the selection requirements. In these cases, we select the combination with the smallest value of $|\cos \theta_{-\pi_f \pi_s} - \cos \theta_{-\pi_f D^*} \cos \theta_{\pi_s D^*}|$. $\theta_{-\pi_f \pi_s}$, $\theta_{-\pi_f D^*}$ and $\theta_{\pi_s D^*}$ are the angles between the directions of $-\vec{p}_{\pi_f}$, \vec{p}_{D^*} and \vec{p}_{π_s} in the $B \rightarrow D^* \pi_f$, $D^* \rightarrow D \pi_s$ decay, respectively. This quantity is zero when the direction of the π_s is a good approximation to the D^* direction and is small when the plane defined by the \vec{p}_{D^*} and \vec{p}_{π_s} nearly coincides with the plane defined by the \vec{p}_{D^*} and $-\vec{p}_{\pi_f}$.

D. Flavor tagging for neutral B meson

The flavor of the signal B decay is obtained from the charge of the fast pion. The flavor of the accompanying B meson decay is determined from the charge of the primary lepton (e or μ) from the semileptonic decay $b \rightarrow c \ell^- \nu$. In addition to tagging the flavor of the B meson decay, there are two other major reasons for using a high-momentum lepton: the point of closest approach of the lepton track and the beam-line provides the location of the tagging side vertex; the requirement of a high momentum lepton in the event dramatically reduces the continuum background.

We require the lepton to have a momentum in the c.m. frame of at least $1.1 \text{ GeV}/c$. We select leptons from well measured tracks by requiring $dr < 0.05 \text{ cm}$, $dz < 2 \text{ cm}$. We demand that the lepton have both $r-\phi$ and z hits in the SVD. To reject secondary leptons from the decay of the unreconstructed D^0 mesons, we require the cosine of the angle between the lepton and π_f to be greater than -0.8 . We also reject lepton tracks that, when combined with any other oppositely charged track in the event, have an invariant mass that is within 50 MeV of the J/ψ mass. If more than one lepton in an event satisfies all of the above criteria, the highest momentum lepton is selected.

However, not all lepton candidates are primary leptons. According to a MC study, the background is dominated by three sources. The first is secondary leptons from charm decays (“cascade leptons”) that come from the $b \rightarrow c \rightarrow s \ell^+ \nu$

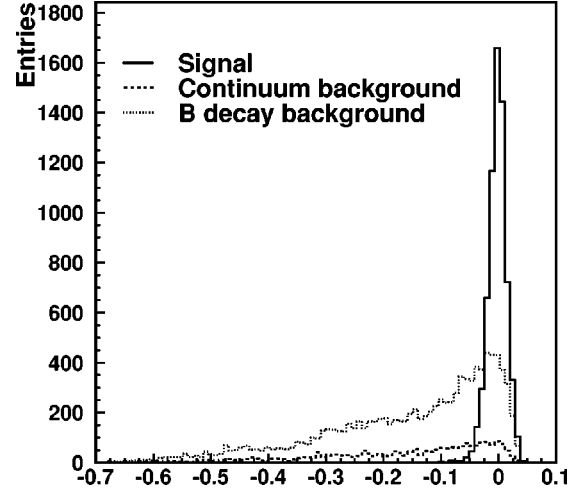


FIG. 2. $\cos \theta_{-\pi_f \pi_s} - \cos \theta_{-\pi_f D^*} \cos \theta_{\pi_s D^*}$ distributions in a simulation of fully reconstructed Monte Carlo events.

decay chain and not directly from $b \rightarrow c \ell^- \bar{\nu}$ like the “primary leptons.” The charges of the cascade and primary leptons are opposite and, thus, cascade leptons can bias the tagged flavor of B mesons and must be well understood. Two of the requirements described above discriminate against secondary leptons: the momentum requirement and the angular cut with respect to the fast π direction. The second category of incorrect tags is due to leptons from J/ψ and $\psi(2S)$ decays. Equal numbers of positively and negatively charged leptons are produced from J/ψ and $\psi(2S)$ decays. Hence, the charge of the observed lepton is uncorrelated with the B flavor. The third category is composed of hadronic tracks misidentified as leptons. Particle identification is applied to suppress this background.

E. Vertex reconstruction

We use the z difference between B decay vertices in the laboratory system to approximate the proper-time difference Δt by

$$\Delta z \equiv z_{\pi_f} - z_l = \gamma(\Delta z_{\text{c.m.}} + c\beta\Delta t) \approx c\beta\gamma\Delta t, \quad (8)$$

where z_{π_f} and z_l are the z positions of the fast pion and lepton tracks respectively, and c is the speed of light. The constant $\beta\gamma = 0.425$ is the Lorentz boost factor for the e^+e^- center of mass system in the Belle experiment. $\Delta z_{\text{c.m.}}$ is the z difference between B decay vertices in the c.m. frame. Here, we assume that the B mesons are at rest in the c.m. frame and, thus, $\Delta z_{\text{c.m.}} \approx 0$. Therefore we have

$$\Delta t \approx \frac{\Delta z}{c\beta\gamma}. \quad (9)$$

The z positions are determined from the intersection of the π_f (or lepton track) with the profile of B decay vertices, which is estimated run-by-run from the profile of the inter-

action point ($\sigma_x^{IP} \sim 110 \mu\text{m}$, $\sigma_y^{IP} \sim 5 \mu\text{m}$, $\sigma_z^{IP} \sim 3500 \mu\text{m}$) convolved with the average B flight length ($30 \mu\text{m}$ in the c.m. frame).

F. Signal

Monte Carlo simulation shows that $B^0(\bar{B}^0) \rightarrow D^{*\mp} \rho^\pm$ candidates are peaked around $1.865 \text{ GeV}/c^2$ in the D^0 missing mass distribution, the same location as the $B^0(\bar{B}^0) \rightarrow D^{*\mp} \pi^\pm$ decays. $B^0(\bar{B}^0) \rightarrow D^{*\mp} \rho^\pm$ decays can also be treated as signal for the measurement of $B^0 - \bar{B}^0$ mixing. Thus, the signal consists of $B^0(\bar{B}^0) \rightarrow D^{*\mp} \pi^\pm$ and $B^0(\bar{B}^0) \rightarrow D^{*\mp} \rho^\pm$ decays tagged by primary leptons. Henceforth, for simplicity, we use the notation $B^0(\bar{B}^0) \rightarrow D^{*\mp} h^\pm$ to represent both $B^0(\bar{B}^0) \rightarrow D^{*\mp} \pi^\pm$ and $B^0(\bar{B}^0) \rightarrow D^{*\mp} \rho^\pm$ decays.

G. Backgrounds

Some other decay modes, such as $B^0 \rightarrow D^{*0} \pi^+$ and $B^0 \rightarrow D^{*0} l^+ \nu_l$, also peak in D^0 missing mass, and hence can fake the signal. We therefore divide the backgrounds into unpeaked and peaked categories. Unpeaked background is dominated by random combinations of π_f and π_s with primary leptons from B^0 and B^\pm decays, and combinatorial background from continuum. We verify in MC that the Δz shape of the unpeaked background can be modeled by the D^0 missing mass sideband. The Δz shapes of unpeaked background thus are taken from the D^0 missing mass sideband. The peaked background is dominated by the following sources: signal B decays [$B^0(\bar{B}^0) \rightarrow D^{*\mp} \pi^\pm$ and $B^0(\bar{B}^0) \rightarrow D^{*\mp} \rho^\pm$] with secondary-lepton tags or fake-lepton tags; $B^0 \rightarrow D^{*0} \pi^+$, $B^+ \rightarrow \bar{D}^{*0} \pi^+$ and $B^0 \rightarrow D^{*0} \pi^+ \pi^0$ decays with primary-lepton tags, secondary-lepton tags or fake-lepton tags. The Δz shapes of the peaked backgrounds are determined from Monte Carlo simulation. The details of the background parametrization are described in Sec. IV.

H. Signal yields

To obtain the signal yields, we apply a binned maximum likelihood fit to the D^0 meson missing mass distribution. In the fit, the signal shape is determined from the $B \rightarrow D^{*0} \pi_f$ Monte Carlo simulation. The shape of the unpeaked background is taken from wrong-sign combinations, where the sign of the charges of the π_f and π_s are the same. The MC simulation shows that the wrong sign combinations have a distribution consistent with the right sign unpeaked background. The peaked background shape comes from the $B\bar{B}$ Monte Carlo simulation. In the fit, the shape of the signal, the peaked background and the unpeaked background are fixed. We also fix the peaked background normalization from the Monte Carlo simulation. We float the normalizations of the signal and unpeaked background. By fitting the D^0 missing mass distributions shown in Figs. 3 and 4, we obtain signal yields of 3433 ± 81 (SF: 751 ± 41 , OF: 2682 ± 70). The estimated backgrounds (unpeaked or peaked) in the sig-

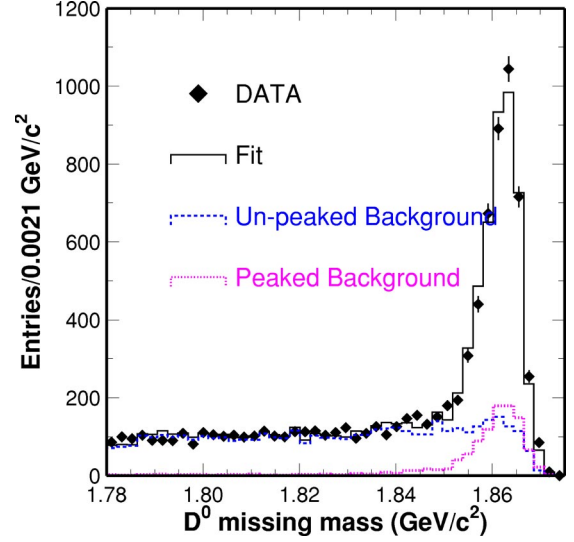


FIG. 3. D^0 missing mass (GeV) distribution for lepton tagged $B^0 \rightarrow D^{*0} h^+$ candidates.

nal region are 1466 ± 37 (SF: $243 \pm 16/219 \pm 15$, OF: $491 \pm 18/513 \pm 23$) for $M_{D_{\text{miss}}} > 1.85 \text{ GeV}/c^2$.

As a consistency check, we use another method to estimate the signal yields. In Fig. 5, we fit the $\cos \theta_{\pi_s}^*$ distribution with both signal and background shapes determined from MC. Here, we float both the signal and background normalizations. For the signal shape, we treat $B^0 \rightarrow D^{*0} \pi^+$ and $B^0 \rightarrow D^{*0} \rho^+$ events separately. The polarization of $B^0 \rightarrow D^{*0} \rho^+$ is fixed from a CLEO measurement [21]. We obtain $B^0 \rightarrow D^{*0} \pi^+$ and $B^0 \rightarrow D^{*0} \rho^+$ yields of 2454 ± 90 and 1353 ± 96 events, respectively. Note that the yields include incorrectly tagged signal. According to a MC study, the probability of incorrect tagging is 9.5%. Thus, we conclude that the two methods yield consistent results. Since the error is smaller for the yields obtained by fitting the D^0 missing mass distribution, we use the results from that method.

IV. MAXIMUM LIKELIHOOD FIT

We extract the mixing frequency Δm_d by simultaneously fitting the time evolution distribution of the SF and OF $B^0 \rightarrow D^{*0} h^+$ samples. The unbinned maximum likelihood fitting method is applied to expressions containing Δm_d as a free parameter, which take into account both signal and background.

A. PDF and likelihood function

Here, we summarize the forms of signal and backgrounds used in the fitting. The likelihood is also established.

For the signal, the probability density functions (PDF) of OF and SF events are given by

$$F_{\text{sig}}^{\text{OF}}(\Delta t) = \int_{-\infty}^{\infty} P^{\text{OF}}(\Delta t') R_{\text{sig}}(\Delta t - \Delta t') d\Delta t', \quad (10)$$

$$F_{\text{sig}}^{\text{SF}}(\Delta t) = \int_{-\infty}^{\infty} P^{\text{SF}}(\Delta t') R_{\text{sig}}(\Delta t - \Delta t') d\Delta t', \quad (11)$$

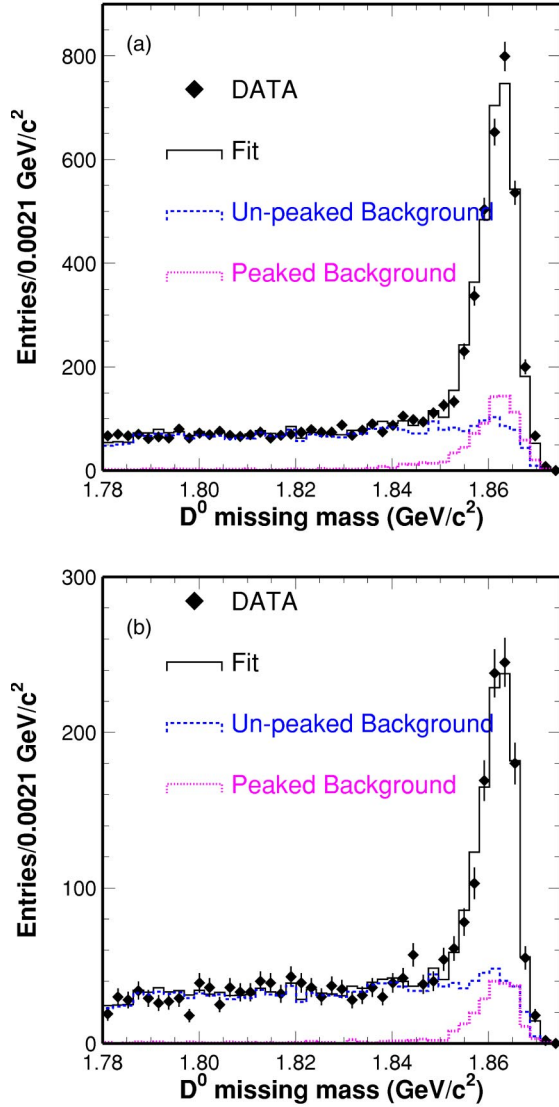


FIG. 4. D^0 missing mass (GeV/c^2) distribution (a) for opposite flavor final states and (b) for the same flavor final states.

where $R_{\text{sig}}(\Delta t)$ is the signal resolution function, which is parametrized by a triple-Gaussian distribution as follows:

$$R_{\text{sig}}(\Delta t) = f_1 G(\Delta t; \mu_1, \sigma_1) + f_2 G(\Delta t; \mu_2, \sigma_2) + (1 - f_1 - f_2) G(\Delta t; \mu_3, \sigma_3). \quad (12)$$

Here, $G(t; \mu, \sigma)$ is the Gaussian distribution, and f_1 and f_2 denote the fractions of the first and second Gaussian, respectively. For the background, the PDFs include separate contributions for the parts that are peaked and unpeaked in the D^0 missing mass distribution. The time dependent parametrization includes prompt (zero lifetime) and finite lifetime components as well as $B^0-\bar{B}^0$ mixing. The time distribution of the unpeaked background is determined from the D^0 missing mass sideband data. The functional forms for the background PDFs are given in the Appendix.

Using the PDFs described in Eqs. (10), (11), (A1) and (A2), the likelihood function can be written as

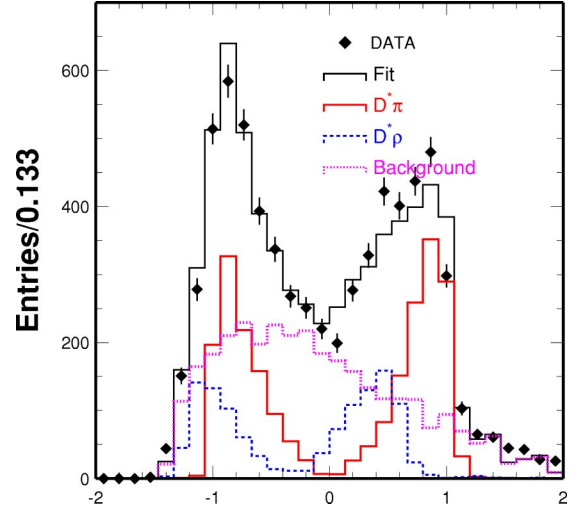


FIG. 5. $\cos \theta_{\pi_s}^*$ distribution. The histogram from the fit and the contributions of signal and background are overlaid.

$$\begin{aligned} \mathcal{L} = & \prod_i ((1 - f_{\text{bkg}}) F_{\text{sig}}^{\text{OF}}(\Delta t_i) + f_{\text{bkg}} f_{\text{bkg}}^{\text{OF}} F_{\text{bkg}}^{\text{OF}}(\Delta t_i)) \\ & \times \prod_j ((1 - f_{\text{bkg}}) F_{\text{sig}}^{\text{SF}}(\Delta t_j) + f_{\text{bkg}} (1 - f_{\text{bkg}}^{\text{OF}}) F_{\text{bkg}}^{\text{SF}}(\Delta t_j)), \end{aligned} \quad (13)$$

where $f_{\text{bkg}} \equiv N_{\text{bkg}} / (N_{\text{bkg}} + N_{\text{sig}})$ is the background fraction. $f_{\text{bkg}}^{\text{OF}}$ is the fraction of OF events in the background and calculated from $N_{\text{bkg}}^{\text{OF}} / N_{\text{bkg}}$. The normalizations are determined from the fit to the D^0 missing mass distribution.

B. Resolution function

The signal resolution function $R_{\text{sig}}(\Delta t)$ is the distribution of the difference between the generated and reconstructed B^0 decay proper times. A good approximation to the resolution function can be obtained from the reconstructed $\Delta z \equiv z_{l^+} - z_{l^-}$ distribution of $J/\psi \rightarrow l^+ l^-$ decays in data, which are parametrized using the triple-Gaussian distribution in Eq. (12). In Monte Carlo simulation, Fig. 6 shows good agreement between the signal resolution function from “ $D^{*-} \pi^+ + l_{\text{tag}}^-$ ” events and the background subtracted Δz distribution from $J/\psi \rightarrow l^+ l^-$ decays. This indicates that using inclusive J/ψ decays to model the tagged $D^* \pi$ resolution function is justified.

$J/\psi \rightarrow l^+ l^-$ candidates are selected using similar requirements to those for the “ $B^0 \rightarrow D^{*-} \pi^+ + l_{\text{tag}}^-$ ” selection. The particle identification criteria are changed to select lepton pairs. Furthermore, no requirement is made on the angle between the tracks, since J/ψ decays produce lepton pairs with large opening angles.² In Fig. 7, the $J/\psi \rightarrow \mu^+ \mu^-$ mass distribution is fitted with the sum of a Gaussian and a second order polynomial. For the $J/\psi \rightarrow e^+ e^-$ case, the sum of a

²Removing this requirement is found not to affect the measured resolution function.

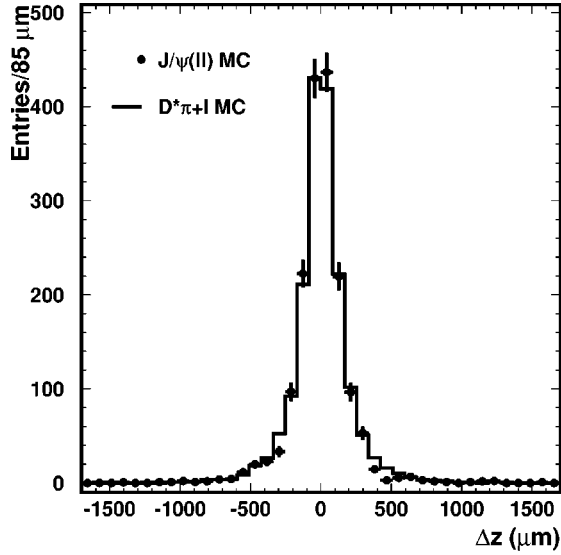


FIG. 6. Resolution functions for $J/\psi \rightarrow l^+ l^-$ decays and “ $D^* \pi^+ + l_{tag}^-$ ” in MC.

“Crystal Ball” function [22] and a second order polynomial is used.

By fitting the J/ψ invariant mass distribution in Fig. 7, we are able to extract the number of signal and background events for both the $e^+ e^-$ and $\mu^+ \mu^-$ cases. To obtain all the parameters of the signal resolution function, we apply a maximum likelihood fit to the overall Δz distribution (see Fig. 8). In the Δz fit, the background shape is determined from the upper sideband of dilepton mass. The normalization is obtained from a second order polynomial fit to the dilepton mass distribution. The background shapes are also parametrized by triple-Gaussian functions. The parameters of the signal resolution function [Eq. (12)] are $\sigma_1 = (94 \pm 6) \mu\text{m}$, $\sigma_2 = (227 \pm 18) \mu\text{m}$, $\sigma_3 = (736 \pm 98) \mu\text{m}$, $f_1 = 0.56 \pm 0.07$ and $f_2 = 0.38 \pm 0.04$. Here, the fit gives mean values for each Gaussian that are consistent with zero.

Similarly, after applying the same technique to MC samples of $J/\psi \rightarrow l^+ l^-$ decays, we find the parameters of the MC resolution function $R_{\text{MC}}(\Delta t)$ to be $\sigma_1 = (79 \pm 4) \mu\text{m}$, $\sigma_2 = (203 \pm 13) \mu\text{m}$, $\sigma_3 = (625 \pm 74) \mu\text{m}$, $f_1 = 0.49 \pm 0.08$ and $f_2 = 0.43 \pm 0.04$. We also fix the mean values to zero.

From the $J/\psi \rightarrow l^+ l^-$ sample (see Fig. 9), we observe a discrepancy between the distributions in data and MC simulation. To account for this difference, we convolve the resolution function in MC with a Gaussian of $50 \mu\text{m}$ width to reproduce the resolution in the data. This additional smearing is applied to the MC simulations of peaked backgrounds. This procedure was verified by comparing the Δz distributions of the D^0 missing mass sideband in data and MC.

For the peaked background case, we use the resolution function for $J/\psi \rightarrow l^+ l^-$ decays in MC with an additional $50 \mu\text{m}$ smearing [$R_{\text{bkg}}^{\text{peak}}(\Delta t) = R_{\text{MC}}(\Delta t) \otimes G(\Delta t; \mu = 0, \sigma = 50 \mu\text{m})$]. With this additional $50 \mu\text{m}$ smearing, we extract the resolution function of the MC simulation. We find that the mean value of this second Gaussian of the resolution function for the same flavor peaked background events is inconsistent with zero, $\mu_2 = (81 \pm 45) \mu\text{m}$. The correspond-

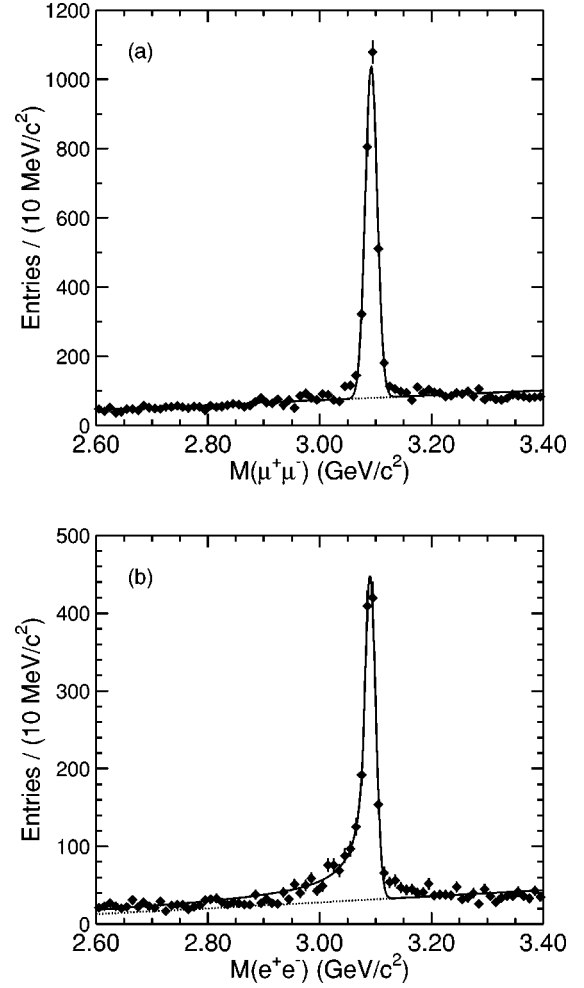


FIG. 7. Invariant mass distributions for (a) $J/\psi \rightarrow \mu^+ \mu^-$ and (b) $J/\psi \rightarrow e^+ e^-$. Fits are described in the text.

ing mean value for the opposite flavor peaked background is consistent with zero. Therefore, an offset is included only in the same flavor peaked background shape.

C. Background distributions

The parametrizations of all background shapes are given in Eqs. (A1) and (A2). Figures 10 and 11 show the Δz distributions for unpeaked and peaked backgrounds, respectively. An unbinned maximum likelihood fitting method is applied to extract the background parameters.

We use the Δz distribution from the D^0 missing mass sideband to reproduce the unpeaked background shape.

For the Δz shapes of the peaked background, we use the Monte Carlo simulation. We find that the probability given by Eq. (A5) for the opposite flavor sample gives a good fit even without the inclusion of the mixing term. The fraction of mixing, f_1 , is then set to zero and fixed for the fit to the opposite flavor sample. The free parameters are f_0 and τ_{bkg} . For the same flavor events, inclusion of the mixing term is necessary to obtain a good fit and the free parameters are f_0 , f_1 and τ_{bkg} . The mixing frequency in the background is determined from the fit.

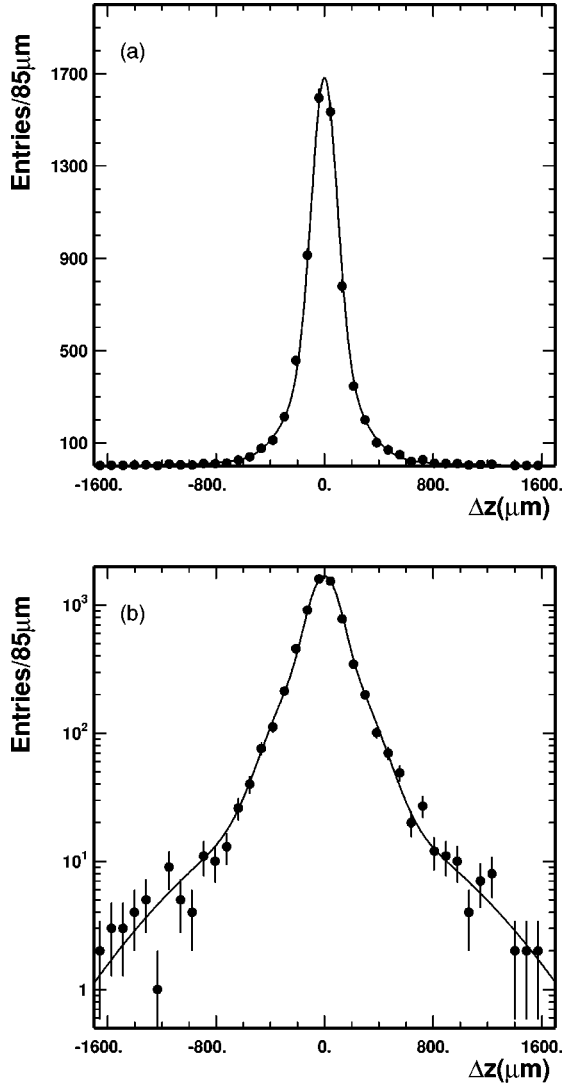


FIG. 8. Signal resolution function fit with a triple-Gaussian for $J/\psi \rightarrow l^+l^-$ decays in the data. Plot (a) shows the fit on a linear scale and plot (b) shows the fit on a semilogarithmic scale.

V. FITTING RESULTS

In the final fit, Δm_d is the only free parameter. The parameters of the resolution functions and background shapes are fixed to their central values. The uncertainties in these parameters are included in the systematic error. The parameters used in the fit are the following:

The B^0 lifetime is fixed at $\tau_{B^0} = (1.548 \pm 0.032)$ ps, the PDG2000 value [19].

The signal and unpeaked background resolution function are determined from $J/\psi \rightarrow l^+l^-$ decays in data. The parameters of Eq. (12) are: $\sigma_1 = (94 \pm 6)$ μm , $\sigma_2 = (227 \pm 18)$ μm , $\sigma_3 = (736 \pm 98)$ μm , $f_1 = 0.56 \pm 0.07$ and $f_2 = 0.38 \pm 0.04$. In the fit, all the parameters are fixed to their central values.

The peaked background resolution function is determined from $J/\psi \rightarrow l^+l^-$ decays in MC convolved with an additional 50 micron smearing term. The parameters of Eq. (12) are: $\sigma_1 = (93 \pm 4)$ μm , $\sigma_2 = (209 \pm 13)$ μm , $\sigma_3 = (627 \pm 74)$ μm , $f_1 = 0.49 \pm 0.08$ and $f_2 = 0.43 \pm 0.04$. In the fit, all the parameters are fixed to their central values.

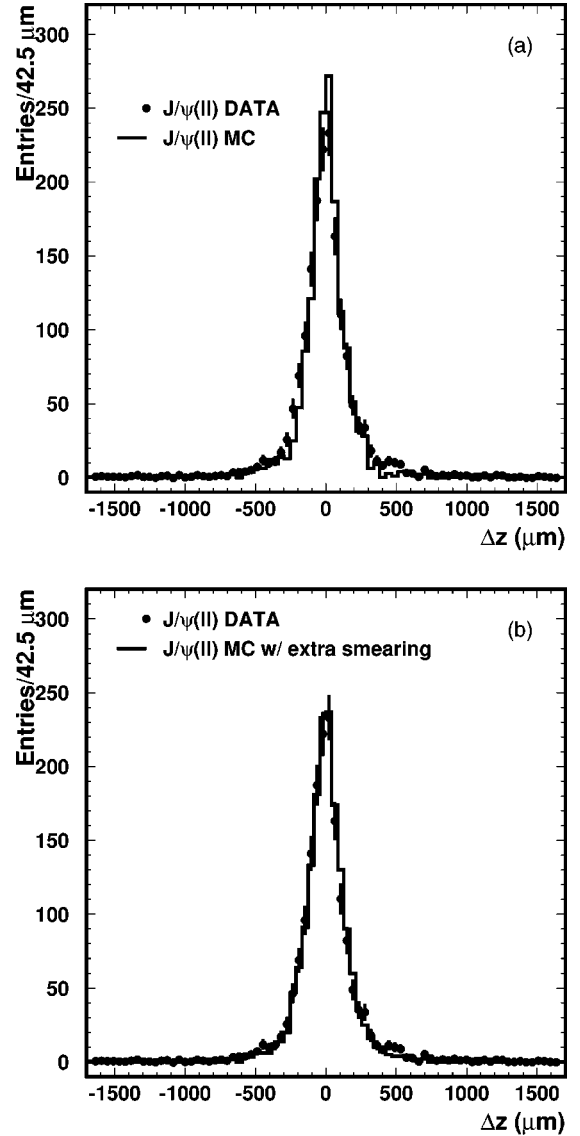


FIG. 9. Δz distribution for $J/\psi \rightarrow l^+l^-$ decays in data and Monte Carlo without (a) and with (b) the extra 50 μm smearing.

Note that the offset of the second Gaussian is $\mu_2 = (81 \pm 45)$ μm for the same flavor peaked background.

The parameters of the unpeaked background shapes are fixed to values determined from fitting the D^0 missing mass sideband (see Fig. 10). The parameters for opposite flavor events are: $f_0 = 0.37 \pm 0.06$, $\tau_{\text{bkg}} = (1.98 \pm 0.17)$ ps. For same flavor events: $f_0 = 0.25 \pm 0.06$, $\tau_{\text{bkg}} = (1.41 \pm 0.10)$ ps.

The parameters of the peaked background shape are fixed to values determined from fitting the Monte Carlo simulation in Fig. 11. Here, Δm_d is fixed to 0.509 ps^{-1} . The parameters for opposite flavor events are: $f_0 = 0.07 \pm 0.07$, $\tau_{\text{bkg}} = (1.34 \pm 0.12)$ ps. The fit gives $f_1 = 0$. For same flavor events: $f_0 = 0.10 \pm 0.04$, $f_1 = 0.40 \pm 0.05$, $\tau_{\text{bkg}} = (1.84 \pm 0.26)$ ps.

The fractions $f_{\text{bkg}} = 0.30 \pm 0.01$, $f_{\text{bkg}}^{\text{OF}} = 0.69 \pm 0.06$, $f_{\text{bkg}}^{\text{OF-unpeak}} = 0.49 \pm 0.03$ and $f_{\text{bkg}}^{\text{SF-unpeak}} = 0.53 \pm 0.04$ are determined from fits to the D^0 missing mass distribution in Fig. 4.

Using the parameters determined above, the Δm_d value extracted from the fit is

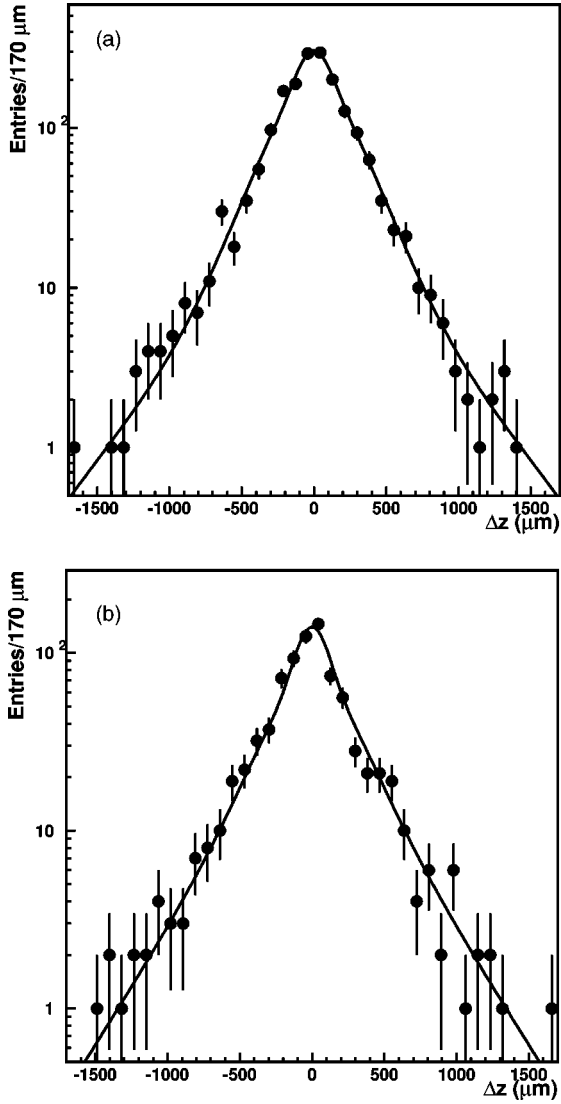


FIG. 10. Semilogarithmic plot of the Δz distributions from the D^0 missing mass sideband for unpeaked backgrounds in (a) opposite-flavor and (b) same-flavor events.

$$\Delta m_d = (0.509 \pm 0.017 \text{ (stat)}) \text{ ps}^{-1}. \quad (14)$$

In Fig. 12 we show the Δz distributions for SF and OF data together with the curves from the fit. To display the charge asymmetry (see Fig. 13) between SF and OF events, we use

$$A(\Delta z) \equiv \frac{N^{\text{OF}}(\Delta z) - N^{\text{SF}}(\Delta z)}{N^{\text{OF}}(\Delta z) + N^{\text{SF}}(\Delta z)}, \quad (15)$$

where $N(\Delta z)$ is the yield of the signal candidates as a function of Δz .

VI. VALIDATION CHECKS

We also perform a series of fits to signal Monte Carlo using the same procedure used in the analysis of the data. We generate four signal MC samples with different Δm_d values. Table I summarizes the fit results and shows the consistency

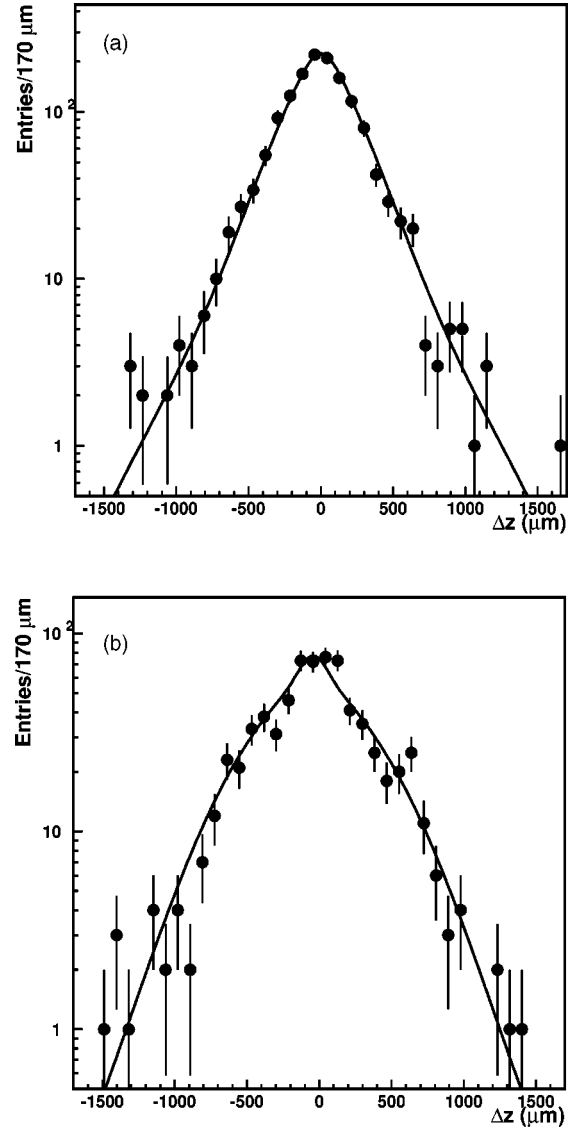


FIG. 11. Semilogarithmic plot of the Δz distributions from MC simulation for peaked backgrounds for (a) opposite-flavor and (b) same-flavor events.

between the input and extracted Δm_d . We conclude there is no significant fit bias.

We also check for run dependence of the Δm_d extraction. The value of Δm_d is shown for four different experimental data taking periods (referred to as experiments 7, 9, 11, 13). The results are shown in Fig. 14(a).

To check the dependence of the Δm_d fit on the continuum background level, we also calculate Δm_d as a function of the cut on the normalized second Fox-Wolfram moment [23] $R_2 = H_2/H_0$. The results are shown in Fig. 14(b).

We have also examined the variation of the fit results on the values of the cuts on the lepton momentum p_l^* in the c.m. frame and the fiducial angle requirement for the lepton and π_f . The results are shown in Figs. 14(c), 14(d). We find that all the variations are consistent with statistical fluctuations.

In addition, we performed fits separately for SF and OF events. The results are shown in Table II.

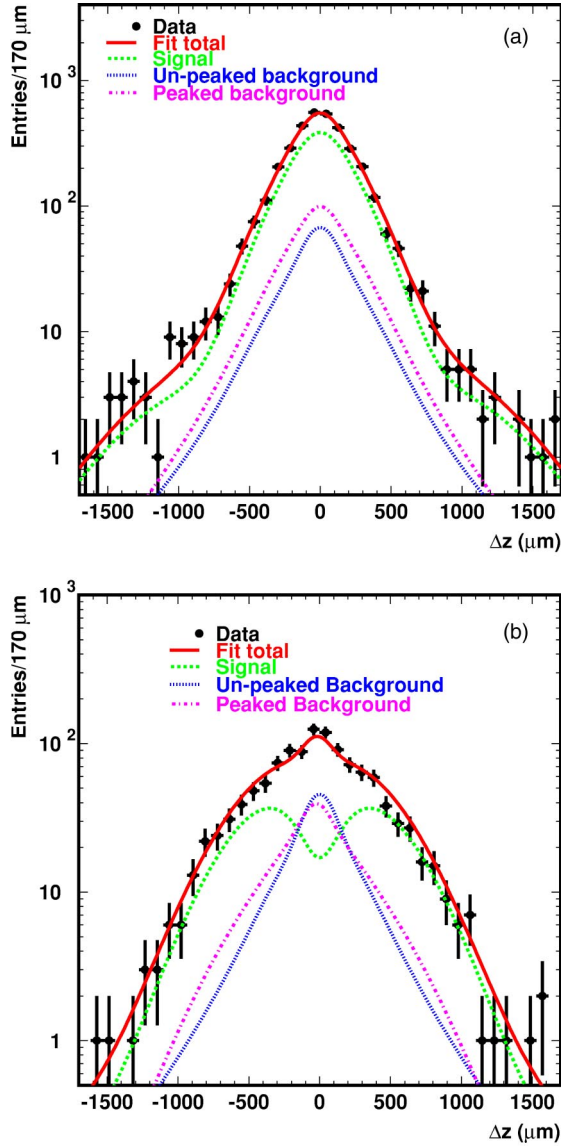


FIG. 12. Δz distributions for (a) opposite-flavor and (b) same-flavor events in data. The curve from the fit and the contributions of signal, unpeaked and peaked backgrounds are overlaid.

The results of separate fits to the e -tagged and μ -tagged events are shown in Table III.

To check the sensitivity of the result to tails of the vertex resolution function, we vary the Δz range of the fit. The results are shown in Table IV and are consistent with the primary result.

VII. SYSTEMATIC UNCERTAINTIES

Various possible sources of systematic uncertainties were investigated. In Table V, we list the contributions to the systematic errors in Δm_d that were estimated by varying the relevant parameters by one standard deviation.

Background fraction

As described in Sec. IV A, the background parts of the likelihood function are weighted by f_{bkg} , $f_{\text{bkg}}^{\text{OF}}$, $f_{\text{bkg}}^{\text{OF-unpeak}}$ and

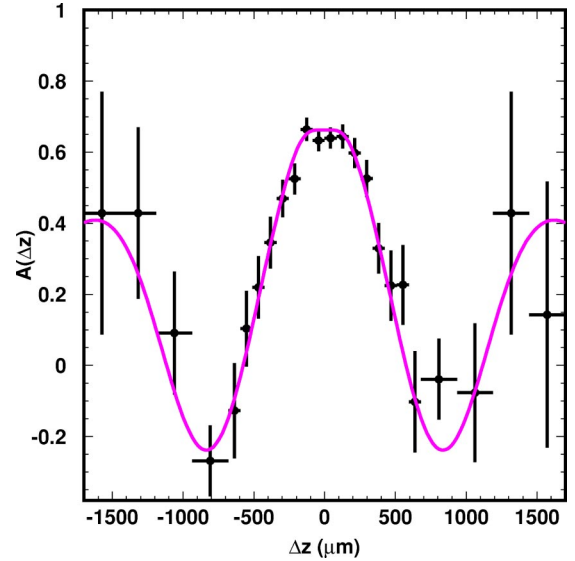


FIG. 13. Distribution of the asymmetry, $A(\Delta z)$, as a function of Δz for data with the fit curve overlaid.

$f_{\text{bkg}}^{\text{SF-unpeak}}$ in Eq. (13); these are determined from the fit to the D^0 missing mass distribution. The corresponding systematic error, ± 0.014 , is estimated by varying the parameters of these background fractions by $\pm 1\sigma$.

Signal resolution function

The fitted parameters of the Δz distributions for $J/\psi \rightarrow l^+l^-$ events are varied according to the limited data statistics by $\pm 1\sigma$. The difference from the standard fit is taken as a measure of the systematic error for each parameter. The corresponding uncertainty for each parameter is added in quadrature to yield a systematic uncertainty in Δm_d of $\pm 0.012 \text{ ps}^{-1}$.

Background shape

For both peaked and unpeaked backgrounds, the Δz distributions are determined from the method described in Sec. IV C. Varying all parameters by $\pm 1\sigma$, we take the differences from the standard fit, ± 0.004 , as the systematic error.

B^0 lifetime

The value of the B^0 lifetime was varied according to the uncertainty of the PDG2000 [19] value. Changes in Δm_d of ± 0.005 are observed.

TABLE I. Summary of signal MC bias test.

Δm_d (ps^{-1}) for MC generation	Δm_d (ps^{-1}) from fit
0.442	0.448 ± 0.011
0.472	0.470 ± 0.012
0.502	0.492 ± 0.011
0.532	0.537 ± 0.012

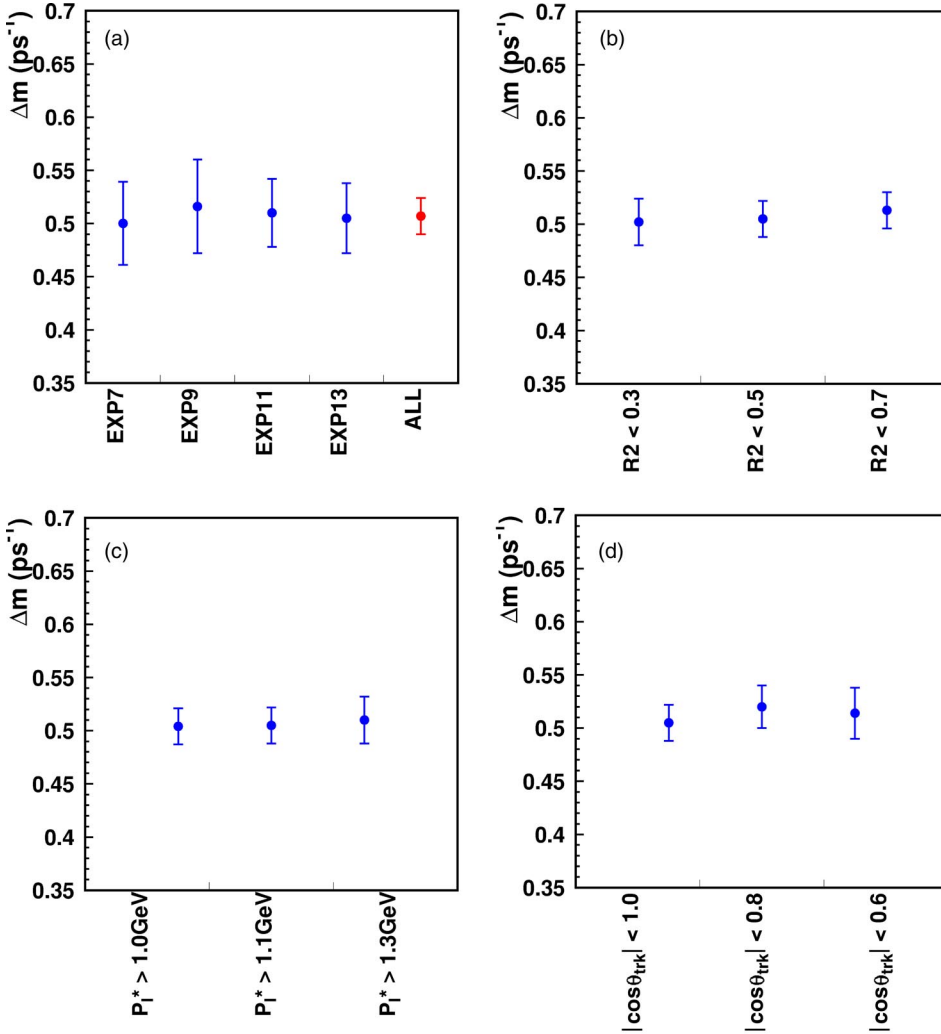


FIG. 14. Δm_d as function of (a) experiment number, (b) R_2 cut value, (c) cut on the lepton momentum in the c.m. frame, and (d) lepton fiducial angle requirement.

Detector resolution

In Sec. IV C, we mentioned that the MC background distributions are smeared to correct for the difference in the vertex resolution between the MC prediction and the data. The smearing function was a Gaussian with $\sigma = 50 \mu\text{m}$. We varied the amount of smearing by its uncertainty ($18 \mu\text{m}$) [24] and repeated the fit. We obtain a difference in Δm_d of ± 0.002 .

The total systematic error, obtained by summing all errors from the different sources in quadrature, is $\pm 0.020 \text{ ps}^{-1}$.

VIII. CONCLUSION

Using 29.1 fb^{-1} of data collected with the Belle detector at the $Y(4S)$, we have measured the $B^0 - \bar{B}^0$ mixing fre-

TABLE II. Summary of Δm_d fit results for SF (OF) events. Only statistical errors are shown.

Sample	Events	Δm_d (ps^{-1})
SF	1213	$\Delta m_d = 0.511 \pm 0.035$
OF	3686	$\Delta m_d = 0.504 \pm 0.019$

quency Δm_d in $B^0(\bar{B}^0) \rightarrow D^{*\mp} \pi^\pm$ decay with a partial reconstruction technique.

The asymmetric e^+e^- beam energies of the KEKB collider allows for the extraction of the time evolution of the B meson wave function from precise measurements of the decay vertex positions. The data are separated into SF and OF samples. The B^0 decay vertex resolution for the signal was determined from the Δz distribution of $J/\psi \rightarrow l^+l^-$ decays that occur in the same data sample. The backgrounds are divided into two components, peaked and unpeaked backgrounds. The Δm_d value, obtained by simultaneously fitting the SF and OF time distributions, is

$$\Delta m_d = (0.509 \pm 0.017(\text{stat}) \pm 0.020(\text{syst})) \text{ ps}^{-1}.$$

TABLE III. Summary of Δm_d fit results for e and μ tagged events. Only statistical errors are shown.

Sample	Events	Δm_d (ps^{-1})
e tagged	2324	$\Delta m_d = 0.510 \pm 0.025$
μ tagged	2575	$\Delta m_d = 0.503 \pm 0.024$

TABLE IV. Summary of Δm_d fit results for different Δz windows. Only statistical errors are shown.

$\max(\Delta z)(\mu\text{m})$	Δm_d (ps $^{-1}$)
1700	0.509 ± 0.017
1275	0.507 ± 0.017
875	0.511 ± 0.019
425	0.520 ± 0.024

This is the first direct measurement of Δm_d using the technique of partial reconstruction of $B^0(\bar{B}^0) \rightarrow D^{*\mp} \pi^\pm$ decays. This measurement is statistically uncorrelated with all other reported experimental results for Δm_d . It is also almost systematically independent from all other measurements. The systematic uncertainty is dominated by the background fractions and the signal resolution function; all quantities are measured experimentally except for the peaked background. This measurement agrees with the world average value of $\Delta m_d = (0.472 \pm 0.017)$ ps $^{-1}$ [19], and serves as a validation of the technique that will be used in the future for the measurement of the CP violation parameter $\sin(2\phi_1 + \phi_3)$ [20].

ACKNOWLEDGMENTS

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APPENDIX: BACKGROUND PDFS

For the background, the PDFs of OF and SF events are parametrized by

TABLE V. Contributions to the systematic error.

Source	Errors (ps $^{-1}$)
Background fraction	± 0.014
Signal resolution function	± 0.012
Background shape	± 0.005
B^0 lifetime (1.548 ± 0.032 ps)	± 0.005
Detector resolution (50 ± 18 μm)	± 0.002
Total	± 0.020

$$F_{\text{bkg}}^{\text{OF}}(\Delta t) = \int_{-\infty}^{\infty} (f_{\text{bkg}}^{\text{OF-unpeak}} P_{\text{bkg}}^{\text{OF-unpeak}}(\Delta t') R_{\text{bkg}}^{\text{unpeak}}(\Delta t - \Delta t') + (1 - f_{\text{bkg}}^{\text{OF-unpeak}}) P_{\text{bkg}}^{\text{OF-peak}}(\Delta t') R_{\text{bkg}}^{\text{peak}}(\Delta t - \Delta t')) d\Delta t', \quad (\text{A1})$$

$$F_{\text{bkg}}^{\text{SF}}(\Delta t) = \int_{-\infty}^{\infty} (f_{\text{bkg}}^{\text{SF-unpeak}} P_{\text{bkg}}^{\text{SF-unpeak}}(\Delta t') R_{\text{bkg}}^{\text{unpeak}}(\Delta t - \Delta t') + (1 - f_{\text{bkg}}^{\text{SF-unpeak}}) P_{\text{bkg}}^{\text{SF-peak}}(\Delta t') R_{\text{bkg}}^{\text{peak}}(\Delta t - \Delta t')) d\Delta t'. \quad (\text{A2})$$

Here, $f_{\text{bkg}}^{\text{SF-unpeak}} \equiv N_{\text{bkg}}^{\text{SF-unpeak}}/N_{\text{bkg}}^{\text{SF}}$ and $f_{\text{bkg}}^{\text{OF-unpeak}} \equiv N_{\text{bkg}}^{\text{OF-unpeak}}/N_{\text{bkg}}^{\text{OF}}$ are the fractions of unpeaked background in the SF and OF events. N is the number of events, the subscripts bkg and sig denote the background and signal components. $R_{\text{bkg}}^{\text{peak}}(\Delta t)$ and $R_{\text{bkg}}^{\text{unpeak}}(\Delta t)$ are resolution functions for peaked and unpeaked background and are parametrized by triple-Gaussian distributions, which have the same form as the signal resolution function $R_{\text{sig}}(\Delta t)$ in Eq. (12). In the fit, we use the signal resolution function for $R_{\text{bkg}}^{\text{unpeak}}(\Delta t)$. The peaked background resolution function $R_{\text{bkg}}^{\text{peak}}(\Delta t)$, described in Sec. IV B, is determined from $J/\psi \rightarrow l^+ l^-$ decays in MC with additional smearing.

In the unpeaked background PDF, $P_{\text{bkg}}^{\text{OF-unpeak}}(\Delta t)$ and $P_{\text{bkg}}^{\text{SF-unpeak}}(\Delta t)$ are defined by

$$P_{\text{bkg}}^{\text{OF-unpeak}}(\Delta t) = f_0^{\text{OF-unpeak}} \delta(\Delta t) + (1 - f_0^{\text{OF-unpeak}}) \frac{1}{2\tau_{\text{bkg}}} e^{-|\Delta t|/\tau_{\text{bkg}}}, \quad (\text{A3})$$

$$P_{\text{bkg}}^{\text{SF-unpeak}}(\Delta t) = f_0^{\text{SF-unpeak}} \delta(\Delta t) + (1 - f_0^{\text{SF-unpeak}}) \frac{1}{2\tau_{\text{bkg}}} e^{-|\Delta t|/\tau_{\text{bkg}}}. \quad (\text{A4})$$

In the peaked background PDF, $P_{\text{bkg}}^{\text{OF-peak}}(\Delta t)$ and $P_{\text{bkg}}^{\text{SF-peak}}(\Delta t)$ are defined by

$$P_{\text{bkg}}^{\text{OF-peak}}(\Delta t) = f_0^{\text{OF-peak}} \delta(\Delta t) + (1 - f_0^{\text{OF-peak}}) \times (1 - f_1^{\text{OF-peak}}) \frac{1}{2\tau_{\text{bkg}}} e^{-|\Delta t|/\tau_{\text{bkg}}} + (1 - f_0^{\text{OF-peak}}) f_1^{\text{OF-peak}} \frac{1 + \Delta m_d^2 \tau_{\text{bkg}}^2}{4\tau_{\text{bkg}} + 2\Delta m_d^2 \tau_{\text{bkg}}^3} \times e^{-|\Delta t|/\tau_{\text{bkg}}} [1 + \cos(\Delta m_d \Delta t)], \quad (\text{A5})$$

$$\begin{aligned}
P_{\text{bkg}}^{\text{SF-peak}}(\Delta t) &= f_0^{\text{SF-peak}} \delta(\Delta t) + (1 - f_0^{\text{SF-peak}})(1 - f_1^{\text{SF-peak}}) \\
&\times \frac{1}{2\tau_{\text{bkg}}} e^{-|\Delta t|/\tau_{\text{bkg}}} + (1 - f_0^{\text{SF-peak}})f_1^{\text{SF-peak}} \\
&\times \frac{1 + \Delta m_d^2 \tau_{\text{bkg}}^2}{2\Delta m_d^2 \tau_{\text{bkg}}^3} e^{-|\Delta t|/\tau_{\text{bkg}}} [1 - \cos(\Delta m_d \Delta t)].
\end{aligned}
\tag{A6}$$

In the above equations; f_0 is the prompt lifetime fraction and $(1-f_0)f_1$ is the mixing fraction; τ_{bkg} is the lifetime of the background component that does not originate from a prompt source; Δm_d contributes to the peaked background, which is dominated by B^0 decays [20]. Mixing terms are included in the PDFs for the peaked background. In the expression for the unpeaked background, we set the mixing term to zero since a parametrization without mixing reproduces the data well.

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- [1] I.I. Bigi and A.I. Sanda, *CP Violation* (Cambridge University Press, Cambridge, England, 2000).
- [2] A.J. Buras *et al.*, Nucl. Phys. **B347**, 491 (1990).
- [3] T. Inami and C.S. Lim, Prog. Theor. Phys. **65**, 297 (1981); **65**, 1772 (1981); A.J. Buras, Phys. Rev. Lett. **46**, 1354 (1981).
- [4] J. F. Donoghue, E. Golowich, and B.R. Holstein, *Dynamics of the Standard Model* (Cambridge University Press, Cambridge, England, 1992).
- [5] I. Dunietz and J.L. Rosner, Phys. Rev. D **34**, 1404 (1986); I. Dunietz, Phys. Lett. B **427**, 179 (1998).
- [6] S. Kurokawa *et al.*, Nucl. Instrum. Methods Phys. Res. A **499**, 1 (2003).
- [7] Belle Collaboration, A. Abashian *et al.*, Nucl. Instrum. Methods Phys. Res. A **479**, 117 (2002).
- [8] G. Alimonti *et al.*, Nucl. Instrum. Methods Phys. Res. A **453**, 71 (2000).
- [9] H. Hirano *et al.*, Nucl. Instrum. Methods Phys. Res. A **455**, 294 (2000); M. Akatsu *et al.*, *ibid.* **454**, 322 (2000).
- [10] H. Kichimi *et al.*, Nucl. Instrum. Methods Phys. Res. A **453**, 315 (2000).
- [11] T. Iijima *et al.*, Nucl. Instrum. Methods Phys. Res. A **453**, 321 (2000).
- [12] H. Ikeda *et al.*, Nucl. Instrum. Methods Phys. Res. A **441**, 401 (2000).
- [13] K. Hanagaki *et al.*, Nucl. Instrum. Methods Phys. Res. A **485**, 490 (2002).
- [14] A. Abashian *et al.*, Nucl. Instrum. Methods Phys. Res. A **449**, 112 (2000).
- [15] The QQ B meson decay event generator was developed by the CLEO Collaboration. See the following URL: <http://www.lns.cornell.edu/public/CLEO/soft/QQ>
- [16] CERN Program Library Long Writeup W5013, CERN, 1993.
- [17] Belle Collaboration, K. Abe *et al.*, Phys. Rev. D **64**, 072001 (2001).
- [18] CLEO Collaboration, R. Giles *et al.*, Phys. Rev. D **30**, 2279 (1984); CLEO Collaboration, G. Brandenburg *et al.*, Phys. Rev. Lett. **80**, 2762 (1998); CLEO Collaboration, B.H. Behrens *et al.*, Phys. Lett. B **490**, 36 (2000).
- [19] Particle Data Group, D.E. Groom *et al.*, Eur. Phys. J. C **15**, 1 (2000).
- [20] Yangheng Zheng, Ph.D. dissertation, University of Hawaii at Manoa, 2002.
- [21] CLEO Collaboration, M.S. Alam *et al.*, Phys. Rev. D **50**, 43 (1994).
- [22] T. Skwarnicki, Ph.D. dissertation, Institute for Nuclear Physics, Krakow, 1986; DESY Internal Report, DESY F31-86-02, 1986.
- [23] G. Fox and S. Wolfram, Phys. Rev. Lett. **41**, 1581 (1978).
- [24] Christos Leonidopoulos, Ph.D. dissertation, Princeton University, 2000.